

### 3 Characteristics of the Bipolar Transistor

The bipolar transistor is a functionally symmetric device—it will work in exactly the same manner if it is hooked up backwards instead of forwards. For most bipolar transistors, this symmetry is only qualitative, because most bipolar transistors are not physically symmetric devices (e.g., typically, the doping density of the emitter is much larger than that of the collector and the cross-sectional area of the emitter-base junction is smaller than that of the collector-base junction).

#### 3.1 The Ebers-Moll Bipolar Transistor Model

The Ebers-Moll model [1] of the bipolar transistor provides a simple, closed-form expression for the collector, base, and emitter currents of a bipolar transistor in terms of its terminal voltages. The model is valid in all regions of operation of the bipolar transistor, transitioning between them smoothly. The Ebers-Moll bipolar transistor model expresses each of the terminal currents in terms of a forward component,  $I_F$ , which only depends on the base-to-emitter voltage and a reverse component,  $I_R$ , which only depends on the base-to-collector voltage. The terminal currents are expressed as

$$I_C = I_F - \frac{I_R}{\alpha_R}, \quad I_E = \frac{I_F}{\alpha_F} - I_R, \quad \text{and} \quad I_B = \frac{I_F}{\beta_F} + \frac{I_R}{\beta_R}, \quad (1)$$

where  $\alpha_F$  and  $\alpha_R$  are, respectively, the forward-active and reverse-active common-base current gains of the bipolar transistor, and  $\beta_F$  and  $\beta_R$  are, respectively, the forward-active and reverse-active common-emitter current gains of the bipolar transistor. These current gains are related to one another by

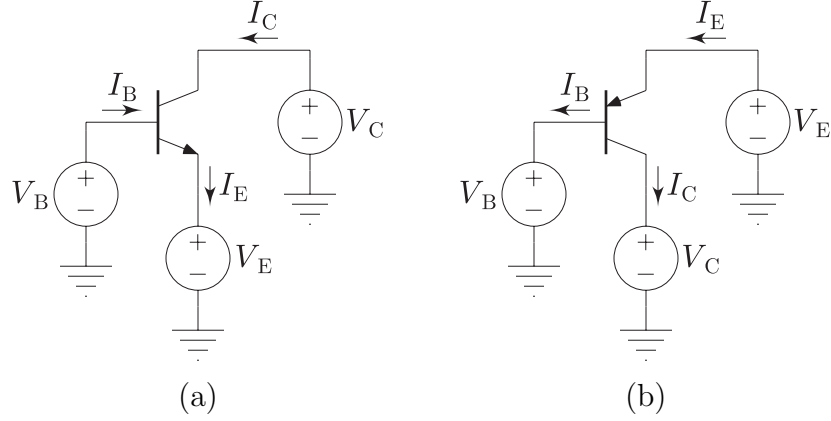
$$\alpha_{F,R} = \frac{\beta_{F,R}}{\beta_{F,R} + 1} \quad \text{and} \quad \beta_{F,R} = \frac{\alpha_{F,R}}{1 - \alpha_{F,R}}.$$

By substituting these relationships into Eq. 1, we can show that the Ebers-Moll model satisfies Kirchhoff's current law, which implies that

$$I_E = I_C + I_B. \quad (2)$$

If the emitter-base junction is forward biased and the collector-base junction is reverse biased, then  $I_F \gg I_R$  and the bipolar transistor is said to be in the *forward-active mode* of operation. In this case, the terminal currents are given by

$$I_C = I_F, \quad I_E = \frac{I_F}{\alpha_F}, \quad \text{and} \quad I_B = \frac{I_F}{\beta_F}. \quad (3)$$



**Figure 1:** Biased bipolar transistors showing the direction of conventional current flow. (a) To bias an *nnp* bipolar transistor into the forward-active region, we should ensure that  $V_{BE} > 4U_T$  and  $V_C \geq V_B$ . In this case,  $I_C$ ,  $I_E$ , and  $I_B$  are positive, flowing in the directions indicated. (b) To bias a *pnp* bipolar transistor into the forward-active region, we should ensure that  $V_{EB} > 4U_T$  and  $V_C \leq V_B$ . In this case,  $I_C$ ,  $I_E$ , and  $I_B$  are positive, flowing in the directions indicated.

By using Eqs. 2 and 3, we can express the terminal currents in terms of one another as

$$I_C = \beta_F I_B = \alpha_F I_E \quad \text{and} \quad I_E = (\beta_F + 1) I_B = \frac{I_C}{\alpha_F}.$$

If the collector-base junction is forward biased and the emitter-base junction is reverse biased, then  $I_R \gg I_F$  and the bipolar transistor is said to be in the *reverse-active mode* of operation, in which the collector and emitters effectively reverse their roles. In this case, the terminal currents are given by

$$I_C = -\frac{I_R}{\alpha_R}, \quad I_E = -I_R, \quad \text{and} \quad I_B = \frac{I_R}{\beta_R}. \quad (4)$$

By using Eqs. 2 and 4, we can express the terminal currents in terms of one another as

$$I_E = -\beta_R I_B = \alpha_R I_C \quad \text{and} \quad I_C = -(\beta_R + 1) I_B = \frac{I_E}{\alpha_R}.$$

If both junctions are forward biased simultaneously, then the bipolar transistor is said to be in *saturation*, which is somewhat unfortunate, as we shall see later in the course, because this operating regime is analogous to the ohmic region of the MOS transistor. The MOS transistor has an operating regime called saturation that is analogous to the active region of the bipolar transistor.

For an *nnp* bipolar transistor, biased as shown in Fig. 1a, the forward current component,  $I_F$ , is given by

$$I_F = I_s \left( e^{V_{BE}/U_T} - 1 \right) = I_s \left( e^{(V_B - V_E)/U_T} - 1 \right), \quad (5)$$

where  $I_s$  is the saturation current of the bipolar transistor and  $U_T$  is the thermal voltage, as before,  $V_B$  is the base voltage, and  $V_E$  is the emitter voltage. Likewise, the reverse current component,  $I_R$ , is given by

$$I_R = I_s \left( e^{V_{BC}/U_T} - 1 \right) = I_s \left( e^{(V_B - V_C)/U_T} - 1 \right), \quad (6)$$

where  $V_C$  is the collector voltage and all of the other symbols are as they were defined previously. By substituting Eqs. 5 and 6 into Eq. 1, we can express the terminal currents in terms of the terminal voltages as

$$\begin{cases} I_C &= I_s \left( e^{(V_B - V_E)/U_T} - 1 \right) - \frac{I_s}{\alpha_R} \left( e^{(V_B - V_C)/U_T} - 1 \right) \\ I_E &= \frac{I_s}{\alpha_R} \left( e^{(V_B - V_E)/U_T} - 1 \right) - I_s \left( e^{(V_B - V_C)/U_T} - 1 \right) \\ I_B &= \frac{I_s}{\beta_F} \left( e^{(V_B - V_E)/U_T} - 1 \right) + \frac{I_s}{\beta_R} \left( e^{(V_B - V_C)/U_T} - 1 \right). \end{cases} \quad (7)$$

If  $V_{BE} > 4U_T$  and  $V_C > V_B$ , then the bipolar transistor is biased into the forward-active region and Eq. 7 becomes

$$I_C \approx I_s e^{(V_B - V_E)/U_T}, \quad I_E \approx \frac{I_s}{\alpha_F} e^{(V_B - V_E)/U_T}, \quad \text{and} \quad I_B \approx \frac{I_s}{\beta_F} e^{(V_B - V_E)/U_T}. \quad (8)$$

If  $V_{BC} > 4U_T$  and  $V_E > V_B$ , then the bipolar transistor is biased into the reverse-active region and Eq. 7 becomes

$$I_C \approx -\frac{I_s}{\alpha_R} e^{(V_B - V_C)/U_T}, \quad I_E \approx -I_s e^{(V_B - V_C)/U_T}, \quad \text{and} \quad I_B \approx \frac{I_s}{\beta_R} e^{(V_B - V_C)/U_T}. \quad (9)$$

If  $V_{BE} > 4U_T$  and  $V_{BC} > 4U_T$ , then the bipolar transistor is biased into the saturation region (or is saturated) and Eq. 7 becomes

$$\begin{cases} I_C &\approx I_s e^{(V_B - V_E)/U_T} - \frac{I_s}{\alpha_R} e^{(V_B - V_C)/U_T} = I_s e^{(V_B - V_E)/U_T} \left( 1 - \frac{1}{\alpha_R} e^{-V_{CE}/U_T} \right) \\ I_E &\approx \frac{I_s}{\alpha_F} e^{(V_B - V_E)/U_T} - I_s e^{(V_B - V_C)/U_T} = \frac{I_s}{\alpha_F} e^{(V_B - V_E)/U_T} \left( 1 - \alpha_F e^{-V_{CE}/U_T} \right) \\ I_B &\approx \frac{I_s}{\beta_F} e^{(V_B - V_E)/U_T} + \frac{I_s}{\beta_R} e^{(V_B - V_C)/U_T} = \frac{I_s}{\beta_F} e^{(V_B - V_E)/U_T} \left( 1 + \frac{\beta_F}{\beta_R} e^{-V_{CE}/U_T} \right). \end{cases} \quad (10)$$

If in addition  $V_{CE}$  is more than a few  $U_T$ , then each of the second terms in parentheses in Eq. 10 becomes negligible compared to one and Eq. 10 further reduces Eq. 8. In this region, which is sometimes called *soft saturation*, the bipolar transistor's behavior is nearly identical to that which it exhibits in the forward-active mode.

If  $V_{BE} < -4U_T$  and  $V_{BC} < -4U_T$ , the bipolar transistor is in the cutoff region and Eq. 7 becomes

$$I_C \approx \frac{I_s}{\beta_R}, \quad I_E \approx \frac{I_s}{\beta_F}, \quad \text{and} \quad I_B \approx -\frac{I_s}{\beta_F} - \frac{I_s}{\beta_R}. \quad (11)$$

For a *pn*p bipolar transistor, biased as shown in Fig. 1b, the forward current component,  $I_F$ , is given by

$$I_F = I_s \left( e^{V_{EB}/U_T} - 1 \right) = I_s \left( e^{(V_E - V_B)/U_T} - 1 \right),$$

where all of the symbols have their previously defined meanings. Likewise, the reverse current component,  $I_R$ , is given by

$$I_R = I_s \left( e^{V_{CB}/U_T} - 1 \right) = I_s \left( e^{(V_C - V_B)/U_T} - 1 \right),$$

where all symbols are as defined previously.

### 3.2 Incremental Bipolar Transistor Characteristics

In this section, we shall develop expressions for and relationships among several small-signal parameters for the bipolar transistor biased in the forward-active region by differentiating the Ebers-Moll model with respect to various terminal voltages. We shall find that restricting our attention to the forward-active region is not really a limitation, because the vast majority of all bipolar integrated circuits make use of these devices in this mode of operation. Moreover, because the bipolar transistor is functionally symmetric with respect to the interchange of its emitter and collector terminals, by using source splitting and superposition, we shall be able to use the incremental relationships that we shall derive in this section for the forward-active bipolar transistor and analogous ones for the reverse-active bipolar transistor to analyze circuits containing bipolar transistors biased into the saturation region.

First, we shall compute the incremental resistance seen looking into the base terminal of an *n*pn bipolar transistor with the emitter voltage held fixed. Recalling that the base current in the forward-active region is given approximately by

$$I_B \approx \frac{I_s}{\beta_F} \cdot e^{(V_B - V_E)/U_T},$$

we can obtain this incremental base resistance by differentiating  $I_B$  with respect to  $V_B$ , which yields

$$\begin{aligned} r_b &= \frac{\partial V_B}{\partial I_B} \\ &= \left( \frac{\partial I_B}{\partial V_B} \right)^{-1} \\ &= \left( \underbrace{\frac{I_s}{\beta_F} \cdot e^{(V_B - V_E)/U_T}}_{I_B} \cdot \frac{1}{U_T} \right)^{-1} \\ &= \frac{U_T}{I_B}. \end{aligned}$$

Next, we shall compute the incremental resistance seen looking into the emitter terminal with the base voltage held constant. Recalling that incremental terminal resistances are

defined on the current flowing into the terminal and that the emitter current in the forward-active region is given approximately by

$$I_E \approx \frac{I_s}{\alpha_F} \cdot e^{(V_B - V_E)/U_T},$$

we can obtain this incremental emitter resistance by differentiating  $-I_E$  with respect to  $V_E$ , which yields

$$\begin{aligned} r_e &= \frac{\partial V_E}{\partial (-I_E)} \\ &= - \left( \frac{\partial I_E}{\partial V_E} \right)^{-1} \\ &= - \left( \underbrace{\frac{I_s}{\alpha_F} \cdot e^{(V_B - V_E)/U_T}}_{I_E} \cdot \left( -\frac{1}{U_T} \right) \right)^{-1} \\ &= \frac{U_T}{I_E}. \end{aligned}$$

Next, we shall define the incremental transconductance gain of the *npn* bipolar transistor as the partial derivative of the collector current with respect to the base voltage. Recalling that the collector current in the forward-active mode of operation is given approximately by

$$I_C \approx I_s e^{(V_B - V_E)/U_T},$$

we can obtain this incremental transconductance gain from

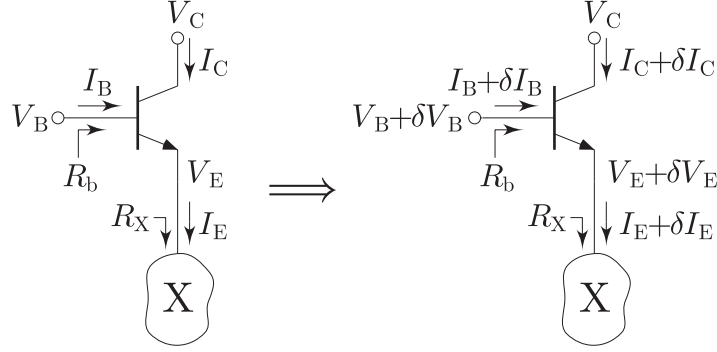
$$\begin{aligned} g_m &= \frac{\partial I_C}{\partial V_B} \\ &= \underbrace{I_s e^{(V_B - V_E)/U_T}}_{I_C} \cdot \frac{1}{U_T} \\ &= \frac{I_C}{U_T}. \end{aligned}$$

Finally, we observe that, because we can express the terminal currents in the forward-active mode in terms of one another according to

$$I_C = \beta_F I_B = \alpha_F I_E, \quad I_E = (\beta_F + 1) I_B = \frac{I_C}{\alpha_F}, \quad \text{and} \quad I_B = \frac{I_C}{\beta_F} = \frac{I_E}{\beta_F + 1},$$

we can also express these small-signal parameters in terms of one another as

$$r_b = \frac{\beta_F}{g_m} = (\beta_F + 1) r_e, \quad r_e = \frac{\alpha_F}{g_m} = \frac{r_b}{\beta_F + 1}, \quad \text{and} \quad g_m = \frac{\beta_F}{r_b} = \frac{\alpha_F}{r_e}.$$



**Figure 2:** Incremental base resistance of an *npn* bipolar transistor with emitter degeneration.

### 3.3 Incremental Base Resistance and Transconductance Gain of an Emitter-Degenerated Bipolar Transistor

In this section, we shall determine what the incremental resistance of the base terminal is when a finite incremental resistance,  $R_X$ , is connected to the emitter terminal, as shown in Fig. 2. Under these circumstances, if we were to increase the base voltage by a small amount, we would increase the base current, the collector current, and the emitter current. However, the increased emitter current must also flow into the circuitry represented by X, which because of its finite incremental resistance,  $R_X$ , the emitter voltage must rise. This increase in the emitter voltage, in turn, reduces the increase in the base, emitter, and collector currents. If less current flows into the base for a given change in the base voltage, we should expect the incremental base resistance under these circumstances should be larger than  $r_b$ . From the incremental base resistance, we shall also determine what effect the emitter degeneration has on the incremental transconductance gain.

For a sufficiently small change in the emitter voltage and current, we have that

$$R_X = \frac{\delta V_E}{\delta I_E},$$

which implies that

$$\delta V_E = \delta I_E \cdot R_X.$$

Because  $I_E = (\beta_F + 1) I_B$ , we can express the change in  $V_E$  in terms of the change in  $I_B$  as

$$\delta V_E = \delta I_B (\beta_F + 1) R_X. \quad (12)$$

Before the small change in the base voltage, we have that the base current is given by

$$I_B = \frac{I_s}{\beta_F} \cdot e^{(V_B - V_E)/U_T}.$$

After the small change in the base voltage and the resulting change in the emitter voltage, we have that

$$I_B + \delta I_B = \frac{I_s}{\beta_F} \cdot e^{((V_B + \delta V_B) - (V_E + \delta V_E))/U_T}$$

$$\begin{aligned}
&= \underbrace{\frac{I_s}{\beta_F} \cdot e^{(V_B - V_E)/U_T}}_{I_B} e^{(\delta V_B - \delta V_E)/U_T} \\
&= I_B \left( 1 + \frac{\delta V_B - \delta V_E}{U_T} + \frac{1}{2} \left( \frac{\delta V_B - \delta V_E}{U_T} \right)^2 + \dots \right) \\
&\approx I_B \left( 1 + \frac{\delta V_B - \delta V_E}{U_T} \right) \\
&= I_B + \frac{\delta V_B - \delta V_E}{r_b},
\end{aligned}$$

which we can rearrange to find that

$$\delta V_B = \delta I_B \cdot r_b + \delta V_E.$$

By substituting Eq. 12 into this equation, we can eliminate  $\delta V_E$ , thereby expressing the change in the base voltage as

$$\delta V_B = \delta I_B \cdot r_b + \delta I_B (\beta_F + 1) R_X,$$

which implies that

$$R_b = \frac{\delta V_B}{\delta I_B} = r_b + (\beta_F + 1) R_X.$$

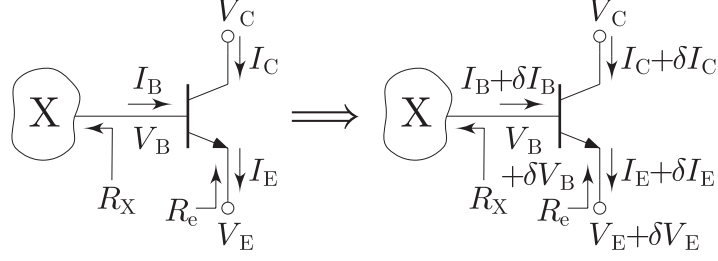
Thus, the incremental base resistance in the presence of emitter degeneration is the unde-generated incremental base resistance in series with the degenerating incremental resistance magnified by  $\beta_F + 1$ .

From this result, we can also obtain the incremental transconductance gain of the bipolar transistor with emitter degeneration. Given a change in  $I_B$ , we can obtain the change in  $I_C$  as

$$\begin{aligned}
\delta I_C &= \beta_F \delta I_B \\
&= \beta_F \cdot \frac{\delta V_B}{R_b} \\
&= \frac{\beta_F \delta V_B}{r_b + (\beta_F + 1) R_X} \\
&= \frac{\beta_F / r_b}{1 + \frac{\beta_F + 1}{r_b} \cdot R_X} \cdot \delta V_B \\
&= \frac{g_m}{1 + R_X / r_e} \cdot \delta V_B,
\end{aligned}$$

which we can rearrange to express the transconductance gain of the bipolar transistor with emitter degeneration as

$$G_m = \frac{\delta I_C}{\delta V_B} = \frac{g_m}{1 + R_X / r_e}.$$



**Figure 3:** Incremental emitter resistance of an *npn* bipolar transistor with a finite incremental resistance in the base.

### 3.4 Incremental Emitter Resistance of an Incremental Base Resistance

In this section, we shall examine the effect on the incremental emitter resistance of a finite incremental resistance connected to the base of a bipolar transistor, as shown in Fig. 3. Given the reference directions for  $I_B$  shown in Fig. 3, for a sufficiently small change in the base voltage and current, we have that

$$R_X = \frac{\delta V_B}{\delta(-I_B)} = -\frac{\delta V_B}{\delta I_B},$$

which implies that

$$\delta V_B = -\delta I_B \cdot R_X.$$

Because  $I_B = I_E / (\beta_F + 1)$ , we can express the change in  $V_B$  in terms of the change in  $I_E$  as

$$\delta V_B = -\delta I_E \cdot \frac{R_X}{\beta_F + 1}. \quad (13)$$

Before the small change in the emitter voltage, we have that the emitter current is given by

$$I_E = \frac{I_s}{\alpha_F} \cdot e^{(V_B - V_E)/U_T}.$$

After the small change in the emitter voltage and the resulting change in the base voltage, we have that

$$\begin{aligned} I_E + \delta I_E &= \frac{I_s}{\alpha_F} \cdot e^{((V_B + \delta V_B) - (V_E + \delta V_E))/U_T} \\ &= \frac{I_s}{\alpha_F} \cdot e^{(V_B - V_E)/U_T} e^{(\delta V_B - \delta V_E)/U_T} \\ &= \underbrace{\frac{I_s}{\alpha_F} \cdot e^{(V_B - V_E)/U_T}}_{I_E} e^{(\delta V_B - \delta V_E)/U_T} \\ &= I_E \left( 1 + \frac{\delta V_B - \delta V_E}{U_T} + \frac{1}{2} \left( \frac{\delta V_B - \delta V_E}{U_T} \right)^2 + \dots \right) \\ &\approx I_E \left( 1 + \frac{\delta V_B - \delta V_E}{U_T} \right) \end{aligned}$$



$$= I_E + \frac{\delta V_B - \delta V_E}{r_e},$$

which we can rearrange to find that

$$-\delta V_E = \delta I_E \cdot r_e - \delta V_B.$$

By substituting Eq. 13 into this equation, we can eliminate  $\delta V_B$ , thereby expressing the change in the emitter voltage as

$$-\delta V_E = \delta I_E \cdot r_e + \delta I_E \cdot \frac{R_X}{\beta_F + 1},$$

which implies that

$$R_e = \frac{\delta V_E}{\delta(-I_E)} = -\frac{\delta V_E}{\delta I_E} = r_e + \frac{R_X}{\beta_F + 1}.$$

Thus, the incremental emitter resistance in the presence of a finite incremental resistance connected to the base is given by  $r_e$  in series with the incremental resistance in the base reduced by a factor of  $\beta_F + 1$ .

## References

- [1] J. J. Ebers and J. L. Moll, “Large-Signal Behavior of Junction Transistors,” *Proceedings of the IRE*, vol. 42, no. ?, pp. 1761–1772, 1954.